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CONTRACTOR REPORT

EXPERIMENTAL INVESTIGATION OF MAGNESIUM

COMBUSTION IN STEAM

Yair Chozev and Jacob Kol

July 1985

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Prepared for: Naval Surface Weapons Center

White Oak Laboratories Silver Springs, MD 20910

NAVAL POSTGRADUATE SCHOOL Monterey, California

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This paper presents a summary of the experimental results and studies of magnesium combustion in steam. The major subject of these studies was to measure the temperature of magnesium particles after the particles had been ejected from an exploding wire. The distribution of particle temperature along the track was measured by the two-color photo-pyrometry method (TC-PPM). In this method, absolute emissivity and velocity of the particles are not

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needed due to differential measurements of radiation energy by density of film exposed to the event. All the other existing methods for moving burning particles require the knowledge of emissivity and velocity correction data. A scanning microdensitometer served as a major data reduction tool in analyzing the film density due to the particle tracks as recorded on calibrated film. The TC-PPH was calibrated using a standard lamp prior to particle temperature measurements.

The initial temperature of the particle ejection from the plasma and at the beginning of combustion is $3430 \text{ K} \pm 9.1\%$. At the initial phase of combustion, the particle consists of pure magnesium with spectral emissivity of about 0.25. During the combustion process, magnesium oxide is formed locally on the particle surface. The oxide is in liquid state with emissivity of 0.9. In the second stage of combustion there are changes on the particle surface in correlation with optical density variation and explained by emissivity variations due to peeling of the oxide film and exposure of the pure magnesium surface. Solidification of the particle, which is magnesium oxide, is observed at the measured temperature of $3125 \text{ K} \pm 5.7\%$. This temperature is near the melting point of magnesia. At the solidification point, abrupt reduction of radiation occurs due to the change of magnesia emissivity from 0.9 (at liquid state) to 0.3 (at solid state).

The spearpoint effect, which is known in the literature as the Blick Effect, was observed at the end of the particle track. The particle supercools below the normal freezing point and then suddenly crystallizes. The measured temperature at the crystallization point is $2840 \text{ K} \pm 5.7\%$.

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ABSTRACT

This paper presents a summary of the experimental results and studies of magnesium combustion in steam. The major subject of these studies was to measure the temperature of magnesium particles after the particles had been ejected from an exploding wire. The distribution of particle temperature along the track was measured by the two-color photo-pyrometry method (TC-PPM), In this method, absolute emissivity and velocity of the particles are not needed due to differential measurements of radiation energy by density of film exposed to the event. All the other existing methods for moving burning particles require the knowledge of emissivity and velocity correction data. A scanning microdensitometer served as a major data reduction tool in analyzing the film density due to the particle tracks as recorded on calibrated film, The TC-PPM was calibrated using a standard lamp prior to particle temperature measurements.

The initial temperature of the particle ejection from the plasma and at the beginning of combustion is 3430 K ± 9.1%. At the initial phase of combustion, the particle consists of pure magnesium with spectral emissivity of about 0.25. During the combustion process, magnesium oxide is formed locally on the particle surface. The oxide is in liquid state with emissivity of 0.9. In the second stage of combustion there are changes on the particle surface in correlation with optical density variation and explained by emissivity variations due to peeling of the oxide film and exposure of the pure magnesium surface. Solidification of the particle, which is magnesium oxide, is observed at the measured temperature of 3125 K ± 5.7%. This temperature is near the melting point of magnesia. At the solidification point, abrupt reduction of radiation occurs due to the change of magnesia emissivity from 0.9 (at liquid state) to 0.3 (at solid state).

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I. INTRODUCTION

The energy released by metals burning in steam has several important applications including torpedo propulsion, nuclear reactor safety, underwater vehicles, underwater ordnance, etc. This report continues the studies that were performed by Hallenbeck [1], Berger et. al., [2] and Kol et. al., [3] at the Naval Postgraduate School which are related to the investigation of underwater shaped charges. Only a few articles dealing with combustion of magnesium in steam have been published; however, none of them were dealing with temperature of magnesium particles burning in steam. The model for magnesium burning in Ar-O₂ mixture that was proposed by Bruzustowski and Glassman [4] and [5] assumed the following characteristics for magnesium combustion:

- (a) At the highest range of burning rates, the flame tearreture will be fixed at the boiling point of the oxide, and the state of the oxide produced will be variable with some oxide always in the condensed state.
- (b) The diffusion of oxygen toward the flame zone is affected by the condensed products of combustion which cannot diffuse and must be convected with the bulk gas motion.
- (c) The existence of condensed species in the flame zone at the high temperature level makes thermal radiation important both in the rate of heat feedback from the flame to the evaporating surface and in the rate of heat loss from the flame to the surroundings.
- (d) The evaporation of metal from the surface may not be very fast in comparison with the diffusion processes occurring in the gas phase.

The degree of dissociation of the products in the flame zone depends upon the balance between heat released by chemical reaction and heat loss by radiation and conduction. For both metal and hydrocarbons the inward

diffusion of oxygen is hindered by the outward diffusion of gaseous products. Grosse and Conway [6] found that the upper limit of the burning temperature of magnesium in oxygen was 3350 K. This result matches the assumption "a" that was made in [4] and [5].

Florko et. al., [7] measured the spectrum of magnesium particle burning in air and the temperature distribution along the distance from the particle center. The temperatures were obtained by using different techniques as follows:

- (a) The surface temperature of the particle was 1400 K measured by using direct thermocouple method.
- (b) The maximum intensity of continuous spectrum temperature was 2600 \pm 200 K at distance of $2r_0$ (r_0 is the particle radius).
- (c) The maximum burning temperature of 2950 ± 200 K was obtained at the distance of $3r_0$ using the relative distribution of the population of the vibrational levels of the excited electronic state assuming Boltzmann distribution.

All of the above experiments were conducted in different conditions compared to experiments reported herein. Therefore one may expect differences in the experimental results. In this report are summarized experimental studies of magnesium burning in steam which included burning temperature and burning behavior.

II. FILM DENSITY VERSUS TEMPERATURE - THEORETICAL MODEL

The exposure of the film, H , ergs/cm² , due to radiation from a burning particle was derived in Berger et. al., [2] as follows:

$$H = \frac{T_{f} \Delta \lambda \varepsilon W_{bb} \pi D}{4f^2 V(t)}$$
 (1)

where

 T_f - Transmissivity of the filters = 0.57 @ 649 nm and 0.58 @ 482 nm

 $\Delta\lambda$ - Bandwidth of the filters = 11.4 nm @ 649 nm and 7.1 nm @ 482 nm

 ϵ - Emissivity of the particle (see equation (2))

 W_{bb} - Black body radiant exitance $\left[\frac{\text{erg}}{\text{cm}^2\text{s}}\right]$

D - Diameter of the particle = 385 µm

f - f-number of the camera (see Table 1)

V(t) - Velocity of the particle* [m/s]

* For the above model initial velocity of 24 m/s was used.

During the combustion process, a liquid oxide film is gradually formed around the pure magnesium; accordingly there is a continuous variation of particle emissivity from pure liquid magnesium emissivity (ε_{m}) to liquid oxide emissivity (ε_{OX}). An emissivity variation is approximated by linear interpolation as follows:

$$\varepsilon_{\mathbf{p}}(\mathbf{T}) = \varepsilon_{\mathbf{m}} \left[1 + \frac{\varepsilon_{\mathbf{O}X} - \varepsilon_{\mathbf{m}}}{\varepsilon_{\mathbf{m}}} \left(\frac{\mathbf{T}_{\mathbf{O}} - \mathbf{T}_{\mathbf{B}}}{\mathbf{T}_{\mathbf{O}} - \mathbf{T}_{\mathbf{B}}} \right) \right]$$
 (2)

where

 ϵ_{OX} - Magnesium emissivity (liquid) = 0.25 @ T < T_B and 0.9 @ T > T_B

 $\varepsilon_{\rm m}$ - Magnesium emissivity (liquid) = 0.25

 ε_{D} - Interpolated particle emissivity

T - Particle temperature during combustion

T_O - Initial particle temperature (3350 K)

 T_R - Temperature at burnout (3075 K)

According to Buber et. al., [8], the emissivity of liquid magnesia above temperature of 3075 K is 0.9. According to Touloukian and DeWitt [9], the emissivity of liquid magnesium is 0.2.

Film optical density at λ_1 = 649 nm can be calculated by the following expression:

$$D_1 = K_1 \log_{10} H_1 - K_1 \log_{10} H_{\min} + D_0$$
 (3)

where

K₁ - Slope of log₁₀H curve at 649 nm

 H_{min} - Minimum exposure energy at gross fog

D_O - Density for gross fog of the film

 H_1 - Film exposure energy from Equation (1) at 649 nm

Accordingly, film optical density at $\lambda_2 = 482$ nm is given by:

$$D_2 = K_2 \log_{10} H_2 - K_2 \log_{10} H_{min} + D_0$$
 (4)

where

K₂ - Slope of log₁₀H curve at 482 mm

H₂ - Film exposure energy from Equation (1) at 482 nm

Using Equations (1) to (4), data for Kodak recording film 2475 [10], and optical instrumentation (see Section III), Figure 1 was obtained. Curves from 1 to 6 of Figure 1 describe film optical density versus temperature for constant emissivity value. Curves from 7 to 12 describe film optical density versus temperature assuming a linear increase of a particle emissivity given by Equation (2). An important observation is that the maximum measured density will be expected near 3200 K. The density will remain high until a temperature of 3075 K (melting temperature of magnesia). Then the density will decay due to the solidification of magnesia.

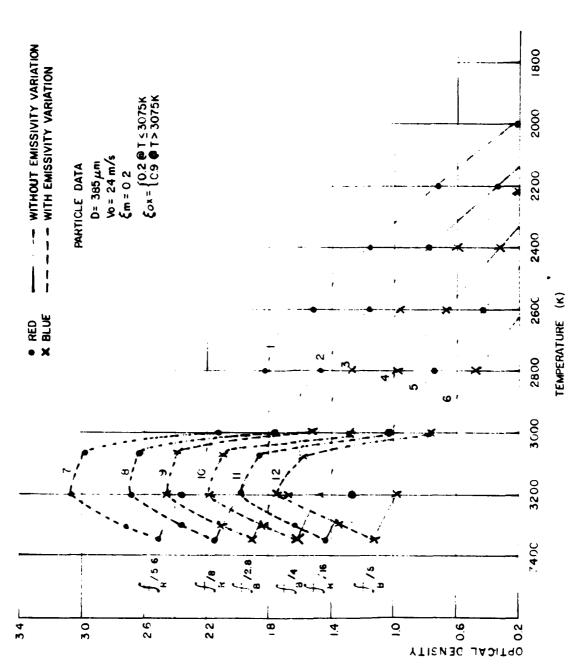


Figure 1. Theoretical Optical Density of the Particle Image Versus Temperature

III. EXPERIMENTAL

The experiments were conducted in pressure vessel which consisted of a twelve inch high stainless steel cylinder, 10.75 inches diameter with four evenly spaced, 5 inch diameter observation ports welded into its circumference. One inch thick, Schlieren quality, Borosicalate crown glass (BK-7) was installed in each port. Two Watlow Band Heaters were used to heat the apparatus to operating temperature and four additional Watlow Heaters were mounted on observation ports in order to prevent steam condensation during experiments. An Omega model 157 Digital Controller was used for temperature stabilization. The experiments were conducted in pressure range of 19 to 32 psi and steam temperature of 180°C. Thermocouples were mounted in different locations inside the chamber to measure the internal temperature.

The magnesium particles were generated by the exploding wire technique.

The 5 cm length wire was mounted between two holders and the ignition energy transferred to the wire was about 90 Joule. The direct energy measurements included the calibrated shunt current measurement and direct voltage measurements across the wire.

Particle temperature was measured by two-color photo-pyrometry method (TC-PPM, see Berger et. al., [2]). An Optronics Microdensitometer Photoscan system P-1000 was used for optical density measurements. Two still Pentax 35 mm cameras were used for two-color (482 nm, 649 nm) photography of the events. Kodak 2475 recording film was calibrated for a detailed graph of film density versus exposure. 1.5 millimeter diameter magnesium wire with purity of 99.998% was used in the experiments. Using the TC-PPM method one can solve Equation (5) using measured densities for λ_1 = 649 nm and λ_2 = 482 nm.

$$T = \frac{C_2(\frac{1}{\lambda_2} - \frac{1}{\lambda_1})}{(\frac{D_1 - D_0}{K_1} - \frac{D_2 - D_0}{K_2}) \text{ on } 10 - \text{ on } A}$$
 (5)

where

$$A = \frac{T_{f1} \quad \lambda_{2}^{5} \quad \Delta \lambda_{1} \quad f_{2}^{2} \quad \epsilon_{1}}{T_{f2} \quad \lambda_{1}^{5} \quad \Delta \lambda_{2} \quad f_{1}^{2} \quad \epsilon_{2}} \stackrel{!}{=} \frac{T_{f1} \quad \lambda_{2}^{5} \quad \Delta \lambda_{1} \quad f_{2}^{2}}{T_{f2} \quad \lambda_{1}^{5} \quad \Delta \lambda_{2} \quad f_{1}^{2}}$$

Assuming that $\frac{\varepsilon_1(649)}{\varepsilon_2(482)}$ = 1 . The assumption that $\varepsilon_1/\varepsilon_2 \approx 1.0$ is important.

See error calculations in [2]. For magnesium particles temperature measurement the following parameters were used:

$$\lambda_1 = 649 \text{ nm} ; \lambda_2 = 482 \text{ nm};$$

$$\Delta\lambda_1 = 11.4 \text{ nm}$$
; $\Delta\lambda_2 = 7.1 \text{ nm}$;

$$T_{f1} = 0.57$$
; $T_{f2} = 0.58$

$$K_1 = 1.2$$
; $K_2 = 0.85$

$$f_1 = 5.6$$
; $f_2 = 2.8$

$$C_2 = 14388 \, \mu \text{m} \, \text{K}$$

Using the above parameters, Equation (5) will be as follows:

$$T = \frac{3338}{\frac{D_1 - D_{01}}{K_1} - \frac{D_2 - D_{02}}{K_2} + 1.056}$$
 (6)

The camera setting and testing conditions for the studies of magnesium particles burning in steam are shown in Table 1.

Table 1. Camera Setting and Testing Conditions for Magnesium Burning in Steam

TEST	AMBIENT	PRESSURE	CAMERA 1			CAMERA 2				
NO.	TEMP.	(PSD)	LENS (mm)	+ NUMBER	λinmi	DISTANCE (cm)*	LENS (mm)	F NUMBER	λınmı	DISTANCE (cm)*
10_	363	21	85	16	649	55	55		482	65
П	360	1919	\$5		649	85_	35	1	482	6.7
12	373	22.5	85		649	\$ 5_	55	1	482	65
13	382	22	15	,	649	55	5.5	4	482	65
14	367	.3()	*5		649	55	5.5	4	482	65
1.5	.383	29	\$	5.6	649	85	7.5	2.5	482	h.i.
16	374	23_	85	7.6	649	5.5	5.5	2.5	482	65
17	1901	23	5.5	5.6	649	55	-55	2.5	182	ьů
18	347	28	85	5.6	649	<u>\$5</u>	55	2.5	482	65
19	377	12	53	5.6	649	85	<u> </u>	2.5	152	65
<u> </u>	360	20	55	` `	649	85	_ 5.5	1	450	65
2	360	20	85	,	649	85	55	1	\$50	65
.3	365	20.5	55	`,	649	55	55	1	450	65
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9	345	20	-55	ī ti_	15.14	-5	- 11	3.5	450	65
[0]	370	20,5	-55	5.6	649	\$5		2.5	450	his
11	360	20	-55	5.6	649	5 5	•	2.5	150	65
12	365	21	*5	535	646	\$5		2.5	\$50	65
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^{*} Distance is the distance of camera 1 and 2 from the source of particles (wire).

IV. RESULTS AND DISCUSSION

A. Photographs of Magnesium Particles Burning in Steam

As discussed previously, the two-color photo-pyrometry method (TC-PPM) requires recording the particle image by calibrated film. Figures 2 and 3 are photographs produced from one exposure in the blue and one in the red wavelength regions.

Figure 2 is the event as photographed with blue filter at 482 nm using an open shutter. Note the ball of plasma formed by the rupture of the magnesium wire. Also note the particle tracks emitting from the plasma. In the original photographs, the spearpoint or Blick Effect can be seen.

Figure 3 is the corresponding photograph for the same event observed with the red filter at 649 nm. The image of Figure 3 is the mirror image of Figure 2. The two cameras observe the exploding wire through observation ports which are 180° apart, i.e., the cameras record the event from opposite directions.

B. Temperature Measurement for Particle Combustion in Steam

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Particle temperature measurements were performed by TC-PPM method which was calibrated prior to the particle combustion experiments by using calibrated quartz-halogen-tungsten lamp with total measured errors of \pm 5.7 percent.

The two film density distributions (in 649 mm and 482 mm wavelengths) along the particle track as function of normalized distance from plasma (distance in experiment = normalized distance x 0.1315 mm) are shown in Figure 4.

The temperature variation along the burning particle track can be seen in Figure 5 and summarized as follows:

The burning particle is ejected from very high temperature plasma. At the

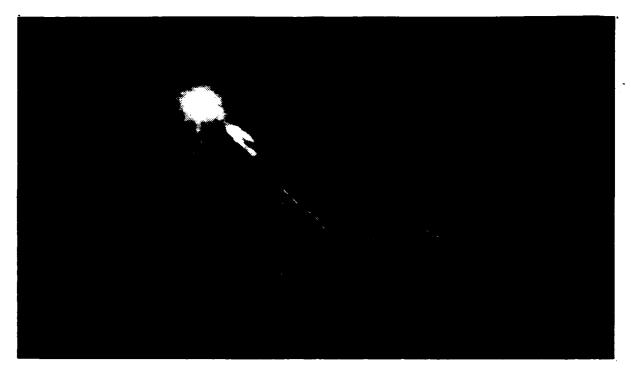


Figure 2. Exploding magnesium wire photographed with blue filter at 482 nm.



Figure 3. Exploding magnesium wire photographed with red filter at 649 nm.

end of the particle track, an abrupt change in measured temperature occurs from 3125 K \pm 5.7% to 1991 K \pm 5.7% as is shown in Figure 5.

At normalized distance less than 120, the temperature calculation was not possible due to high intensity radiation at λ_2 = 482 nm. Further discussion concerning this phenomenon will be given at next paragraph.

C. Burning Behavior of Magnesium Particles in Steam

Five different tracks of film density variation are shown in Figure 4 and in the Appendix from Figures 11a to 11d. According to the film density variation of a representative track in Figure 4, one can distinguish three different burning regions as follows:

The first region, for normalized distance less than about 120 is characterized by D_2 (blue) greater or equal to D_1 (red). The second region, for normalized distance from 120 to 210 (approximately) is characterized by continuously growing difference between the red and the blue densities $(D_1 - D_2)$. In the third region, for normalized distance greater than 210, the difference between densities is decreasing.

- a. First Region Blue Radiation Interference
 - In order to explain the source of the blue radiation interference in the first region, the following studies were performed:
 - 1. By examination of the microdensitometer results for this region, higher blue densities were observed at longer radial distances from the plasma center than at red densities. It is shown in Figure 6.
 - 2. Magnesium wires burning in steam were photographed by 35 mm color film, and a blue glowing "cloud" was observed around the plasma with streaks propagating outwards.

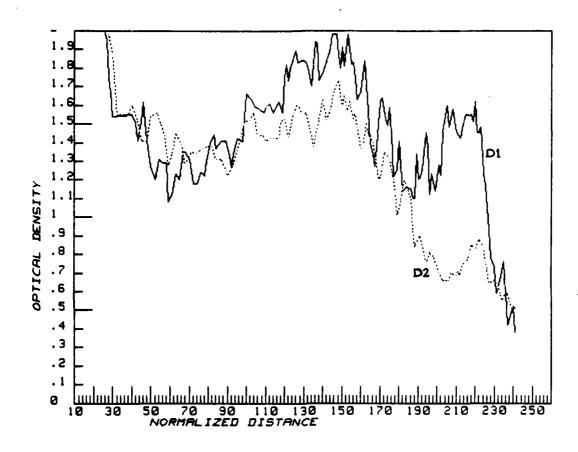


Figure 4. Measured film density as a function of normalized distance along a typical particle track.

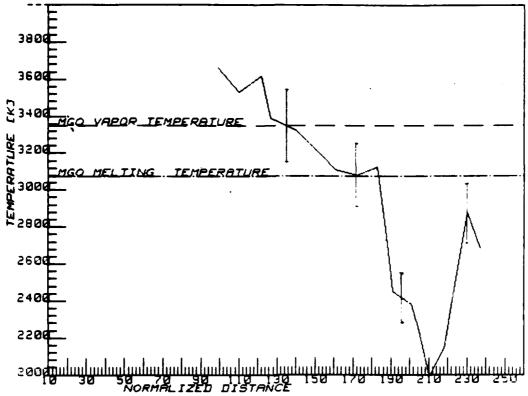


Figure 5. Typical temperature variations of magnesium particle in steam along the particle track (based on the densities of Figure 4).

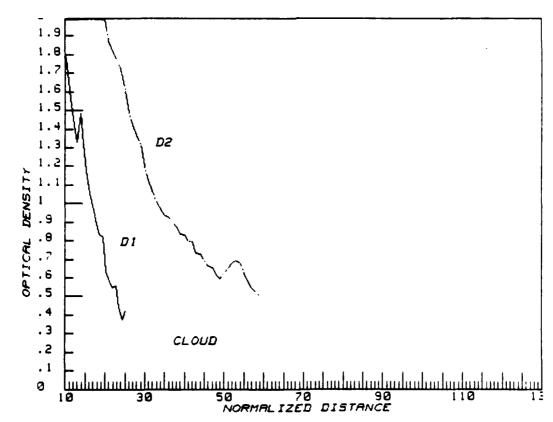


Figure 6. Optical density due to cloud of particles versus distance along particle track.

- 3. The spectroscopic studies that were performed by Brzustowski and Glassman [5] showed that when magnesium ribbons are burning in an ambient gas consisting of 40% oxygen and 60% argon at a pressure of 300 mm Hg, the reaction emits Mg* and MgO* spectral lines. The most relevant lines are the excited MgO peak lines from 491 to 501 nm. Other strong lines were the green lines of Mg* (~517 nm).
- 4. The possibility of emission from other various components is summarized using [11], [12] and [13] in Table 2. As is shown there is no evidence of other spectral lines such as of MgOH and O₂.
- 5. In order to investigate the source of the blue emission around the plasma there were performed experiments of igniting magnesium wires in two different pressures of air: 0.25 and 14 psi. The 35 mm color photographs showed that in the higher pressure of air, the blue cloud was about the same size as in steam, considering the difference of pressures. In the lower pressure atmosphere, the blue "cloud" was significantly smaller, but there was an interface of green radiation between the white plasma and the blue (smaller) "cloud". The behavior of ignited magnesium wires is shown in Figure 7. The green radiation indicates magnesium vapor (see (3) above) around the plasma in the low pressure air.

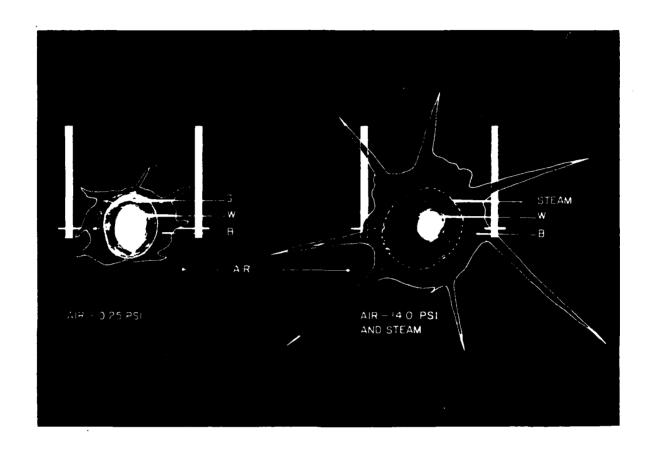


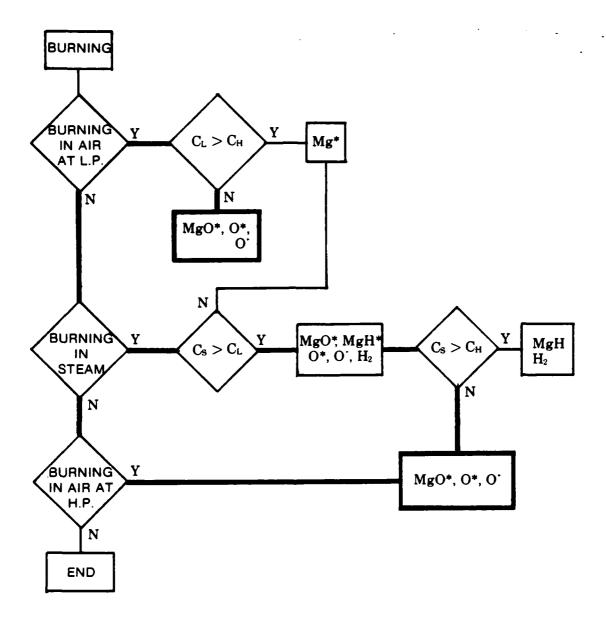
Figure 7. Plasma description in air and in steam.

- G Green
- B Blue
- W White

Table 2. Spectral Lines of Mg, O, H, H2, MgH and MgO from 470 mm to 490 mm.

		RELATIVE	
COMPOUND	λ(nm)	INTENSITY	eV
Mg	470.3	7	7.0
Mg+	474.0	5	14.2
Mg+	485.1	7	14.2
0	480.3	4	13.3
	477 • 4	5	13.3
	477.3	4	13.3
υ +	489.1	4	28.8
	487.2	5	31.4
	485.7	3	28.8
	475.1	4	28.8
	471.0	5	28.8
Н	(H _B)486.1	500 (?)	12.7
H_2	487.3	8-10	
_	485.7	8-10	i
	484.9	8-10	İ
!	482.3	8-10	Ī
	476.4	8-10	j
MgH	484.5		j
Mg0	481.9	3	ĺ
J	481.0	3	İ
	480.1	3	

- 6. By using the logical flow graph in Figure 8 and the condensed data from 1 to 5, one can conclude that the designated path in Figure 8 is the most reasonable for magnesium burning in steam, and the possible sources of blue radiation are MgO*, O* and O*. By comparing the burning of Mg to other metals burning in steam such as Al, Ta, Zr and pyrofuze no evidence of blue radiation was observed around the plasma; therefore, the choice of O* and O* is not relevant for magnesium. In addition, the excitation energies of O* and O* (13.3 eV, 30 eV respectively) are much higher then excitation energy of MgO (2.5 eV); therefore we believe that excited MgO is the source of the blue radiation as was found by Brzustowsky and Glassman in air [5].
- 7. The interference of the "cloud" of the blue radiation was measured by microdensitometer up to a normalized distance of 60 (Figure 6), but the interference along the tracks is up to normalized distance of 120. One possible reason for this behavior is as follows: The "cloud" (smoke) of blue radiation consists of small particles distributed along the track (10 µm and less [5]), which are cooling faster than the measured big particles (100 µm 400 µm) [1] therefore one can conloude that excited MgO exists along particle track up to normalized distance of 120 due possibly to high initial heat transfer from plasma. The particle gradually transfers the heat to the ambient gas and finally, at the end of first region, excited MgO does not exist and MgO occurs in the ground state.



 $C_{\rm L}$ - Diameter of the blue "cloud" in air at low pressure

 $C_{\mbox{\scriptsize H}}$ - Diameter of the blue "cloud" in air at high pressure

 C_{S} - Diameter of the blue "cloud" in steam at high pressure

Figure 8. Logical flow graph for magnesium experiments in steam (20 psi) and air at two different pressures (0.25, 14 psi).

Second Region - Changes in Magnesium Particle Surface This region begins at normalized distance of 120 and ends at normalized distance of 210 (see Figure 4). There is observed increase of measured density of the film as was photographed with blue filter at 482 nm and with red filter at 649 nm. The increase of the density from normalized distance of 120 to 150 could be related to an increase of emissivity due to successive covering of the particle by oxide. This behavior also was obtained by the theoretical studies as was described in Figure 1. Photographs 2 and 3 show changes on the particle surface, and there is correlation between these changes and the rapid changes in the measured densities of the film in Figure 4. Derevyaga et. al., [14] suggest the following explanation to those changes: The droplet begins to lose the oxide film; a zone of vapor gas reaction forms about it; and the brightness of the radiation from the particle increases. By correlating this oxide peeling to emissivity variations, while the variation of the temperature is only 305 K, one can calculate the change in emissivity by using equation (6) that was developed in [2]:

$$\frac{\varepsilon_{\text{ox}}}{\varepsilon_{\text{m}}} = \log_{10}^{-1} \frac{D_{\text{max}} - D_{\text{min}}}{K_{i}}$$
 (6)

where D_{max} - The maximum density of the film

 D_{min} - The minimum density of the film

K_i - The slope of the logH curve

By using the measured results of the film densities, one can obtain the ratio of emissivities $\epsilon_{\rm OX}/\epsilon_{\rm m}$ as follows:

at 482 nm the ratio was from 1.7 to 2

at 649 nm the ratio was from 1.81 to 2.28

Assuming that the emissivity at the steady state of magnesia at temperature above 3075 K is 0.9 (Buber et. al., [8]) the average change in emissivity will be from 0.9 to 0.486 at 482 nm and from 0.9 to 0.44 at 649 nm. Those results are comparable to the results that were published by Touloukian and DeWitt [9] for liquid magnesium ($\varepsilon_{\rm m}$ = 0.2 - 0.25) considering the assumption that the oxide peeling is not complete.

c. Third Region - Near The Spearpoint

Along the track, in normalized distance greater than 210, the radiation from the particle becomes less intense, then brightens, gradually loses intensity again and then disappears. An explanation of this phenomena based on super cooling of the molten oxide particle has been discussed by Nelson et. al., [15]. According to this mechanism, when the burning is near completion, the particle temperature is below the normal freezing point, i.e., the particle is in a supercooled condition. Then the particle suddenly crystallizes and increases to a temperature near but below the equilibrium melting temperature with an abrupt increase of thermal emission. Finally, the particle cools gradually to ambient temperature.

The abrupt increase of temperature to 2840 K \pm 5.7% was measured by TC-PPM for magnesium burning particle in steam. The amount of temperature increase, ΔT , can be calculated from heat of solidification as follows:

$$\Delta T = \frac{\Delta H_f}{C_p} = 1350 \text{ K}$$

where ΔH_f , heat of solidification, 459 cal/gm, C_p , heat capacity of magnesia at 2500 K, 0.34 cal/gmK.

The calculated amount of temperature increase and the measured amount of temperature increase of $1130 \pm 5.7\%$ are comparable.

D. Temperature Investigation by Using Different Filters

In order to eliminate the influence of the radiation in 482 nm, there were performed a series of experiments in the following wavelengths:

- 1) 450 nm and 649 nm
- 2) 599 nm and 649 nm.

By using the first set of filters in the experiments (Table 1) the influence of spectral line radiation was still significant. By using the second set of filters in other experiments (Table 1) there were measured optical densities by using a microdensitometer and are shown in Figure 9. Calculated temperatures are shown in Figure 10 by using the following measured parameters of the 599 nm filter:

 $T_{f} = 0.535$

 $\lambda = 599 \text{ mm}$

 $\Delta \lambda = 10.5 \text{ nm}$

Therefore one can conclude that the temperature for magnesium burning in steam in first region is close to the MgO vapor temperature. The total error for temperature measurement (see for details [2]) in this set of filters is 9.1% and is shown in Figure 10.

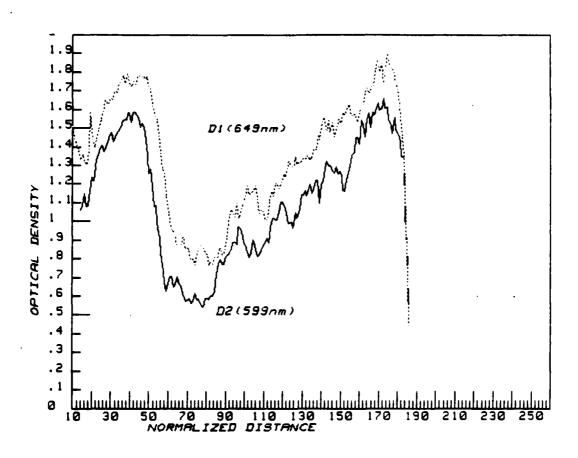


Figure 9. Measured film density in wavelength of 649 nm and 599 nm as a function of normalized distance along the particle track.

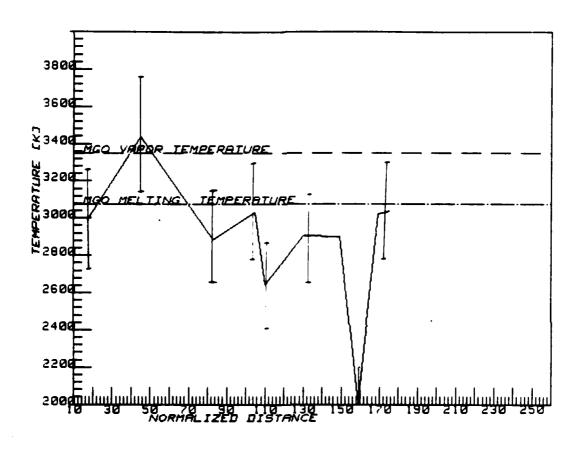


Figure 10. Typical temperature variations of magnesium particle in steam along the particle track based on the densities of Figure 9.

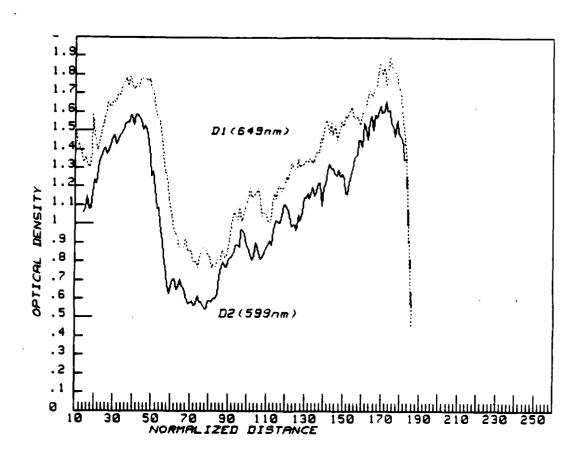


Figure 9. Measured film density in wavelength of 649 nm and 599 nm as a function of normalized distance along the particle track.

V. SUMMARY

The temperature measurement of magnesium particles burning in steam by using TC-PPM method can be summarized as follows: The maximum temperature of the burning particles along the track was about the boiling temperature of magnesia (3350 K). By observing the photgraphic results and the measured densities of the film, one can observe three regions of different behavior along the particle tracks as follows:

a. The First Region

At close range to the plasma the calculated temperature based on experimentally measured densities was about the vapor temperature of magnesia (3350 K). The radiation in this region is more intense in wavelength ranges from 482 to 450 nm compared to 649 nm. This phenomenon were found by comparing the results of measured optical densities by using different filters.

b. The Second Region

In this region the influence of the blue radiation which was measured at 482 nm and 450 nm, is reduced and the cooling of the particle was observed. The temperature decreased from 3430 K \pm 9.1% to 3125 K \pm 5.7%, while the temperature is decreased by only small amount, the changes in radiation (measured from density) in this region are relatively high; see Figure 4. The changes could be related to emissivity variations due to peeling of the magnesia as was seen also in Figures 2 and 3.

c. The Third Region

This region is related to the end of burning. The particle temperature is 1990 K and is below the normal freezing point of magnesia (3075 K), i.e., the magnesia is supercooled. Then suddenly

the magnesia crystallizes and returns to near the equilibrium melting temperature with abrupt increase of thermal emission that was measured as 2840 K \pm 5.7%; see Figure 5.

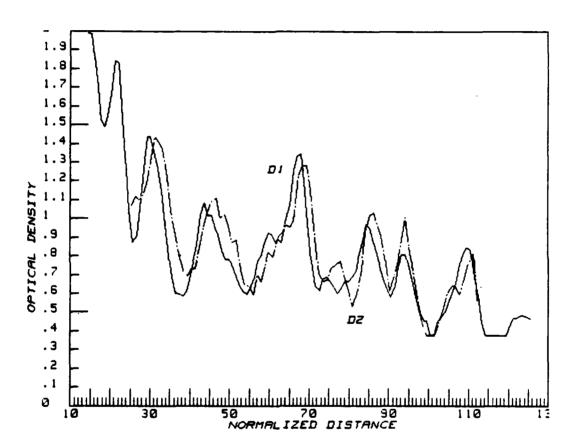


Figure 11. Measured film density as a function of normalized distance along the particle track.

a. Track 2

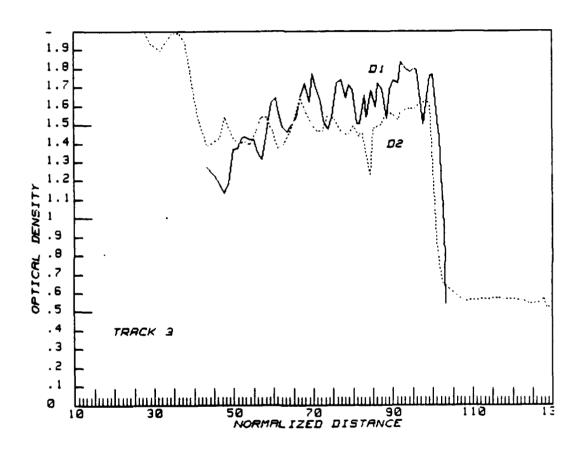


Figure 11. Measured film density as a function of normalized distance along the particle track.

b. Track 3

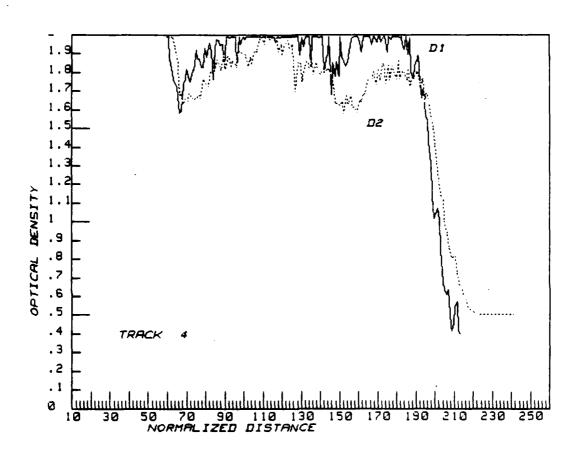


Figure 11. Measured film density as a function of normalized distance along the particle track.

c. Track 4

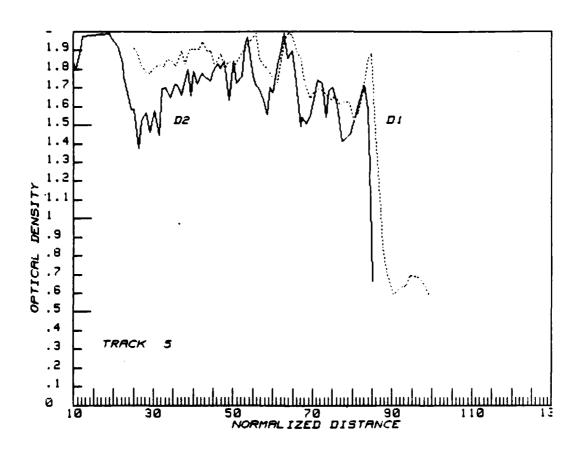


Figure 11. Measured film density as a function of normalized distance along the particle track.

d. Track 5

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